

# **Cosmic Tsunamis in Modified Gravity**

## scalar waves wrecking screening mechanisms

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With R. Hagala and C. Llinares



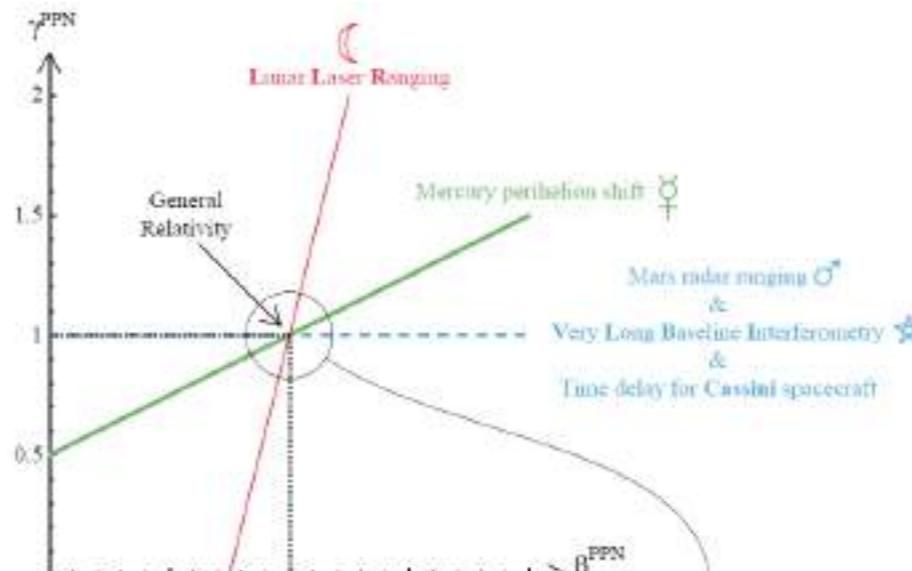
UiO : Institute of Theoretical Astrophysics  
University of Oslo

# Extending General Relativity

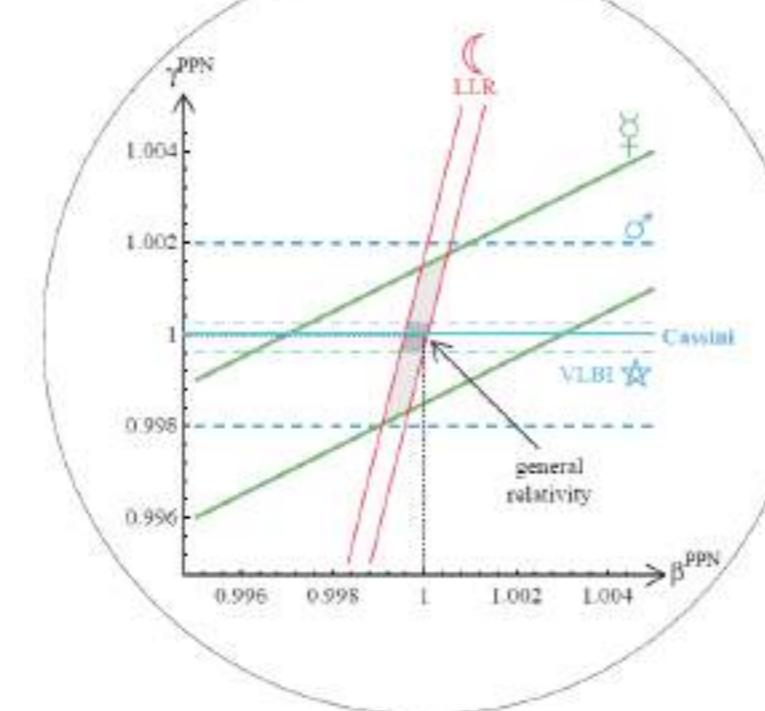
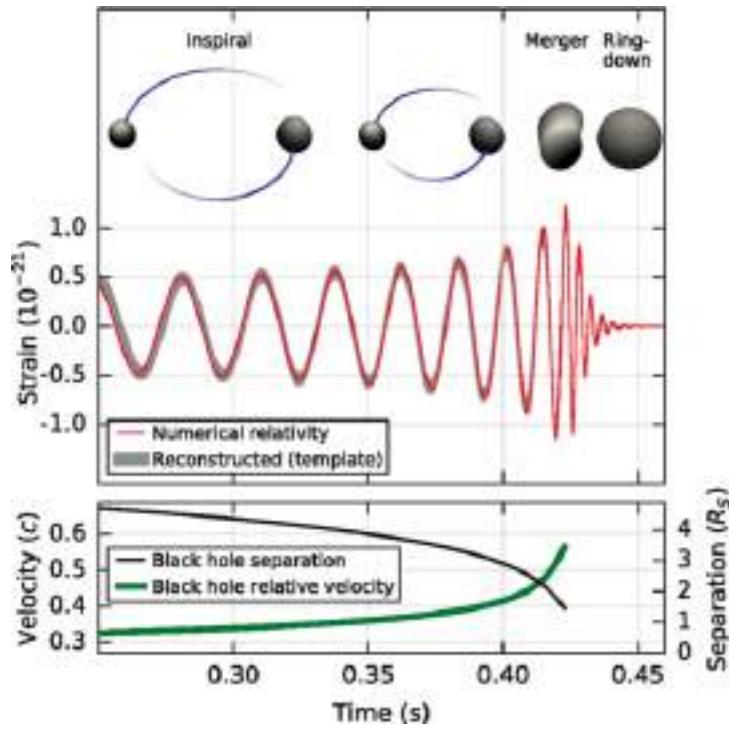
Light degree of freedom driving acceleration

## Screening Mechanisms!

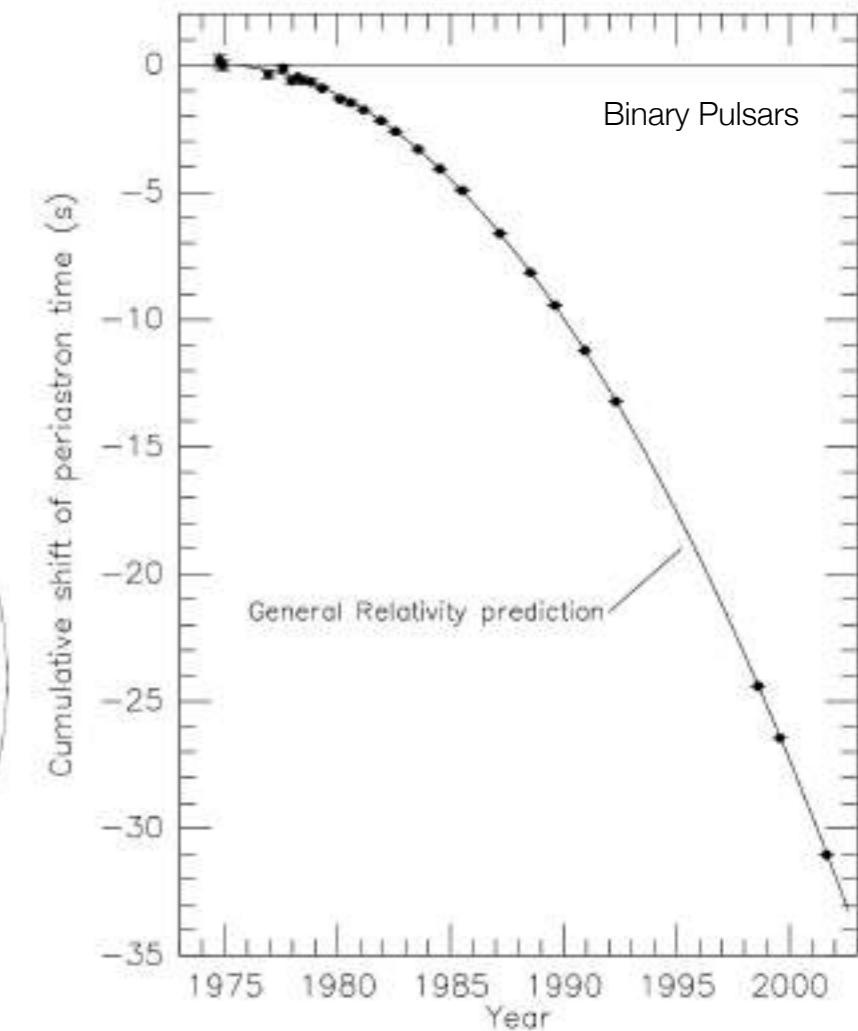
### Solar System Bounds



### Strong field Bounds



### Astrophysical Bounds



# *Screening Mechanisms in Scalar-Tensor Gravity*

## *Suppressing the Yukawa Potential*

$$\mathcal{L} = -\frac{1}{2} Z^{\mu\nu}(\phi, \partial\phi, \dots) \partial_\mu \phi \partial_\nu \phi - V(\phi) + g(\phi) T_\mu^\mu$$

**1. Weak coupling**

**2. Large mass**

**3. Large inertia**

$$V(r) = -\frac{g^2(\bar{\phi})}{Z(\bar{\phi})c_s^2(\bar{\phi})} \frac{e^{-\frac{m(\bar{\phi})}{\sqrt{Z(\bar{\phi})c_s(\bar{\phi})}} r}}{4\pi r} \mathcal{M}$$

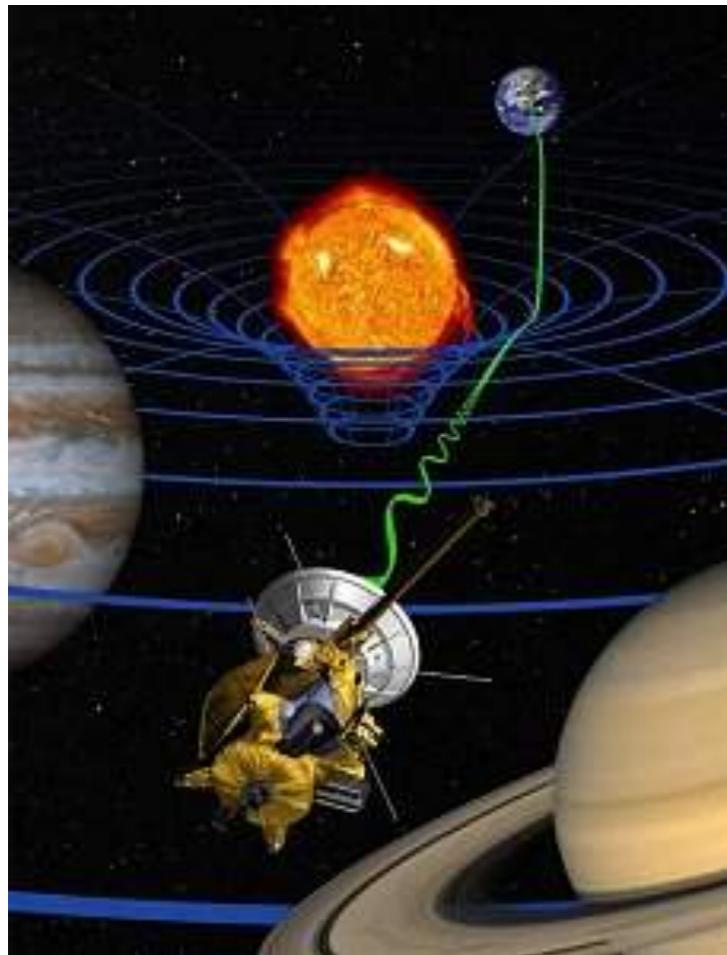
*Efficiency crucial to evade gravity constraints*

# Constraints on Modified Gravity

## Post-Newtonian Parameter

$$ds^2 = \left(-1 + 2\frac{GM}{r}\right) dt^2 + \left(1 + 2\gamma\frac{GM}{r}\right) dx^2$$

“How much space is curved by a unit rest mass”



$$\gamma - 1 = (2.1 \pm 2.3) \times 10^{-5}$$

$$\gamma - 1 = 0, \quad \text{GR}$$

$$\gamma - 1 = -\frac{\phi^2}{M^2} \frac{2}{\frac{\phi^2}{M^2} + 2\Psi\left(1 + \frac{\phi^2}{M^2}\right)}$$

# Computing the profile of the field in the solar system

$$S = \int d^4x \sqrt{-g} \left( \frac{M_{Pl}^2}{2} R - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right) + S_{\text{matter}} [A^2(\phi) g_{\mu\nu}, \psi]$$

- Scalar field equation of motion

$$\ddot{\phi} + 3H\dot{\phi} - \frac{1}{a^2} \nabla^2 \phi = -V_{\text{eff},\phi}(\rho, \phi)$$

- A damped **wave equation**

## Quasi-static approximation

Field profile does not change in virialised/quasi-static systems

- Scalar field equation of motion

$$\frac{1}{a^2} \nabla^2 \phi = -V_{\text{eff},\phi}(\rho, \phi)$$

- **Quasi-static** approximation

$$S = \int d^4x \sqrt{-g} \left( \frac{M_{\text{Pl}}^2}{2} R - \frac{1}{2} (\partial\phi)^2 - V(\phi) \right) + S_{\text{matter}} [A^2(\phi) g_{\mu\nu}, \psi]$$

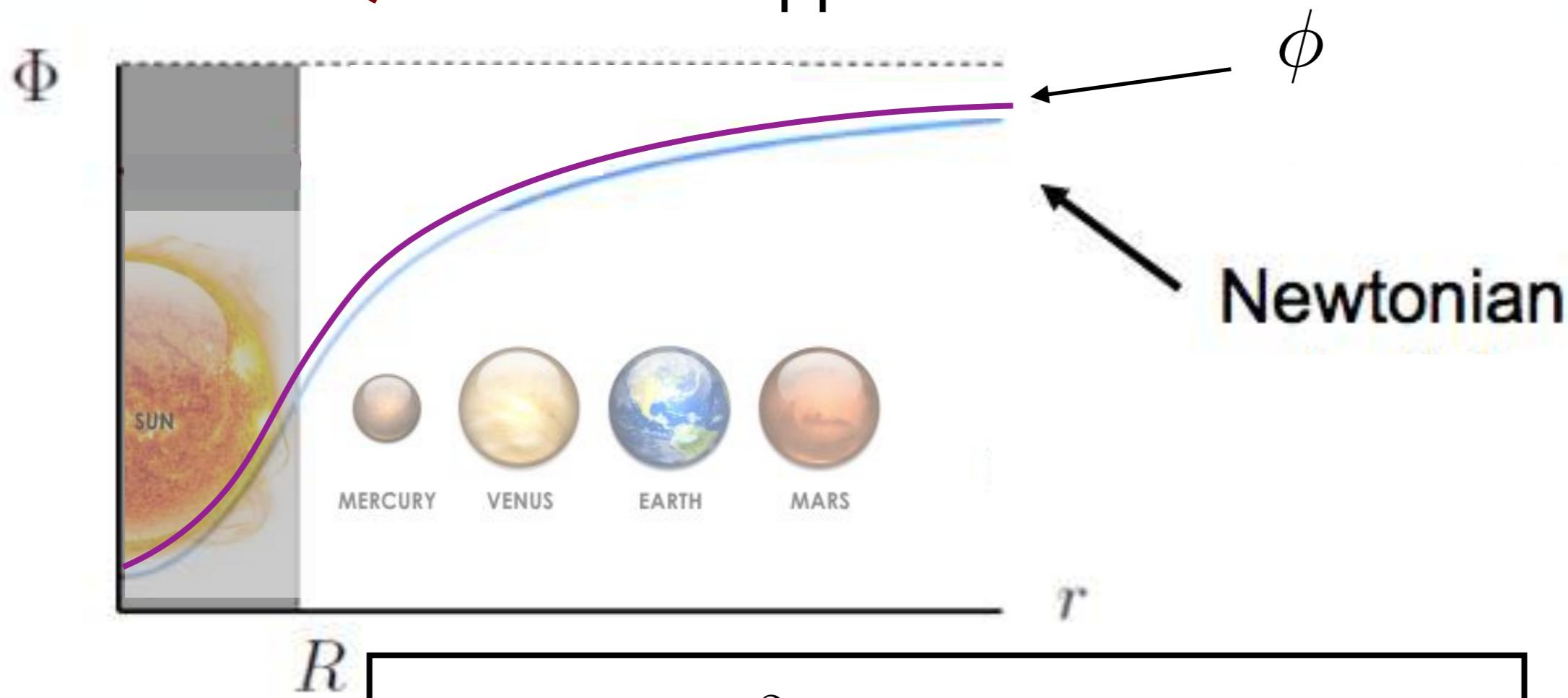
# Quasi-static approximation

Field profile does not change in virialised/quasi-static systems

- Scalar field equation of motion

$$\frac{1}{a^2} \nabla^2 \phi = -V_{\text{eff},\phi}(\rho, \phi)$$

- **Quasi-static** approximation

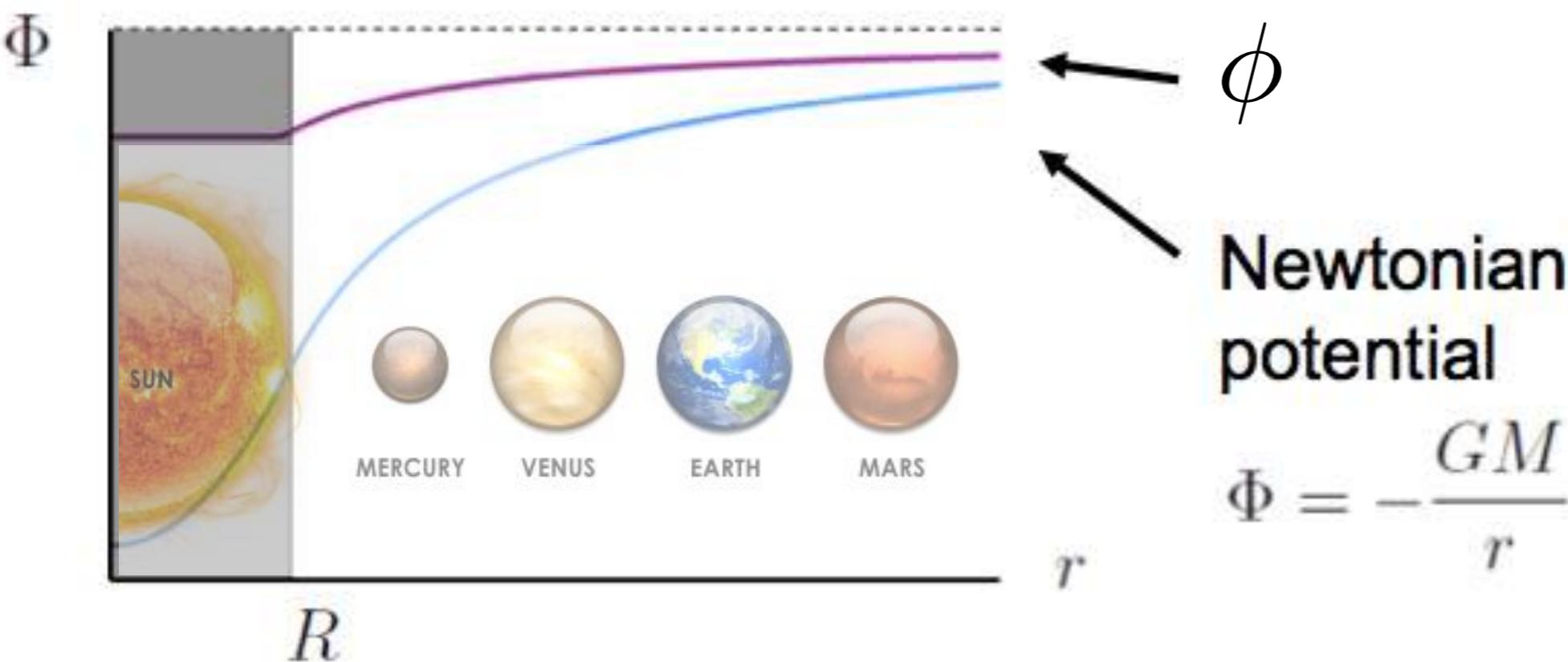


$$\gamma - 1 = -\frac{\phi^2}{M^2} \frac{2}{\frac{\phi^2}{M^2} + 2\Psi(1 + \frac{\phi^2}{M^2})} \ll 1$$

# Screening mechanism suppress field value and gradient

$$\vec{a} = -\vec{\nabla}\Phi - \frac{d \ln A(\phi)}{d\phi} \vec{\nabla}\phi$$

$$F_\phi \ll F_N$$



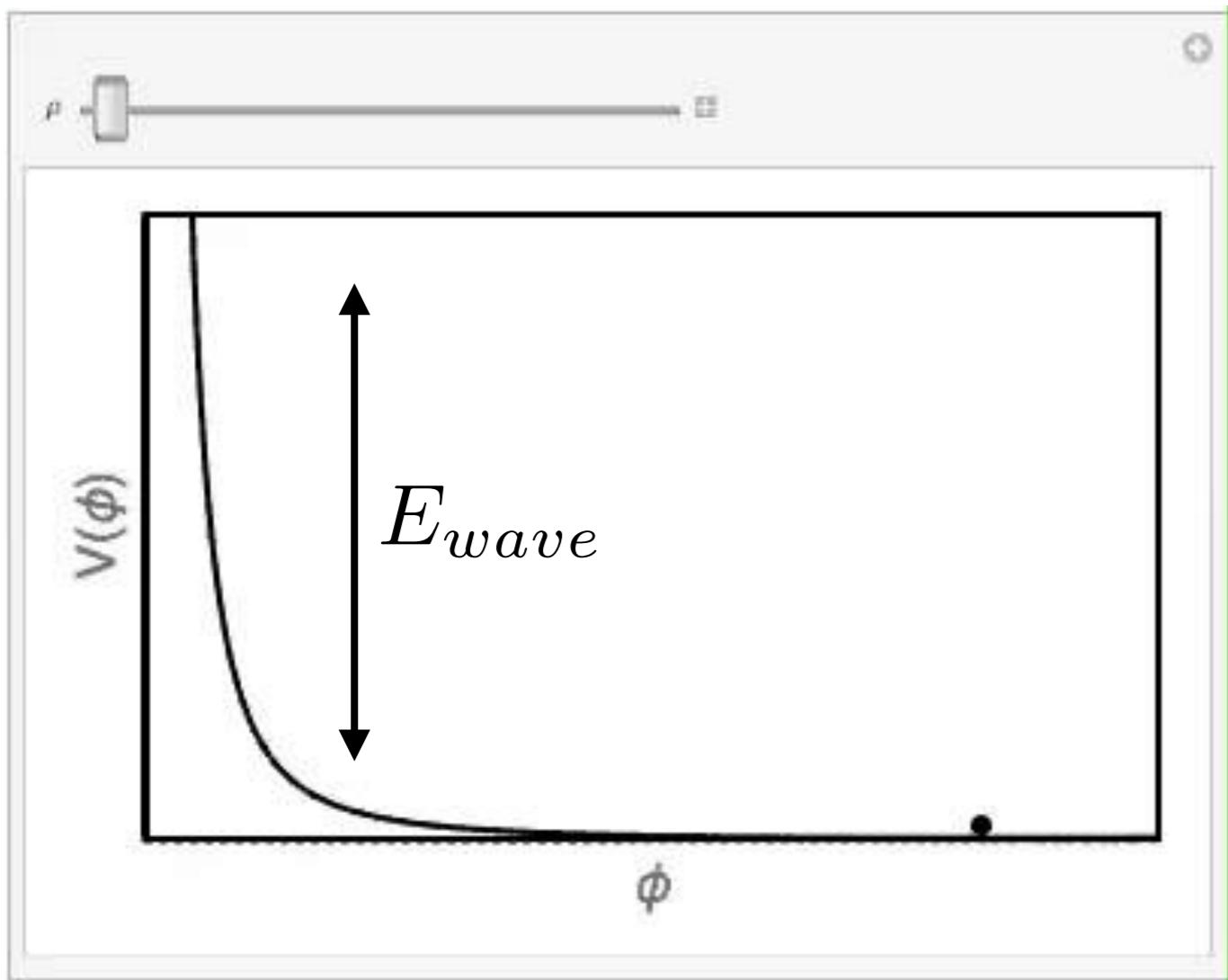
$$\Phi = -\frac{GM}{r}$$

$$\gamma - 1 = -\frac{\phi^2}{M^2} \frac{2}{\frac{\phi^2}{M^2} + 2\Psi(1 + \frac{\phi^2}{M^2})} \ll 1$$

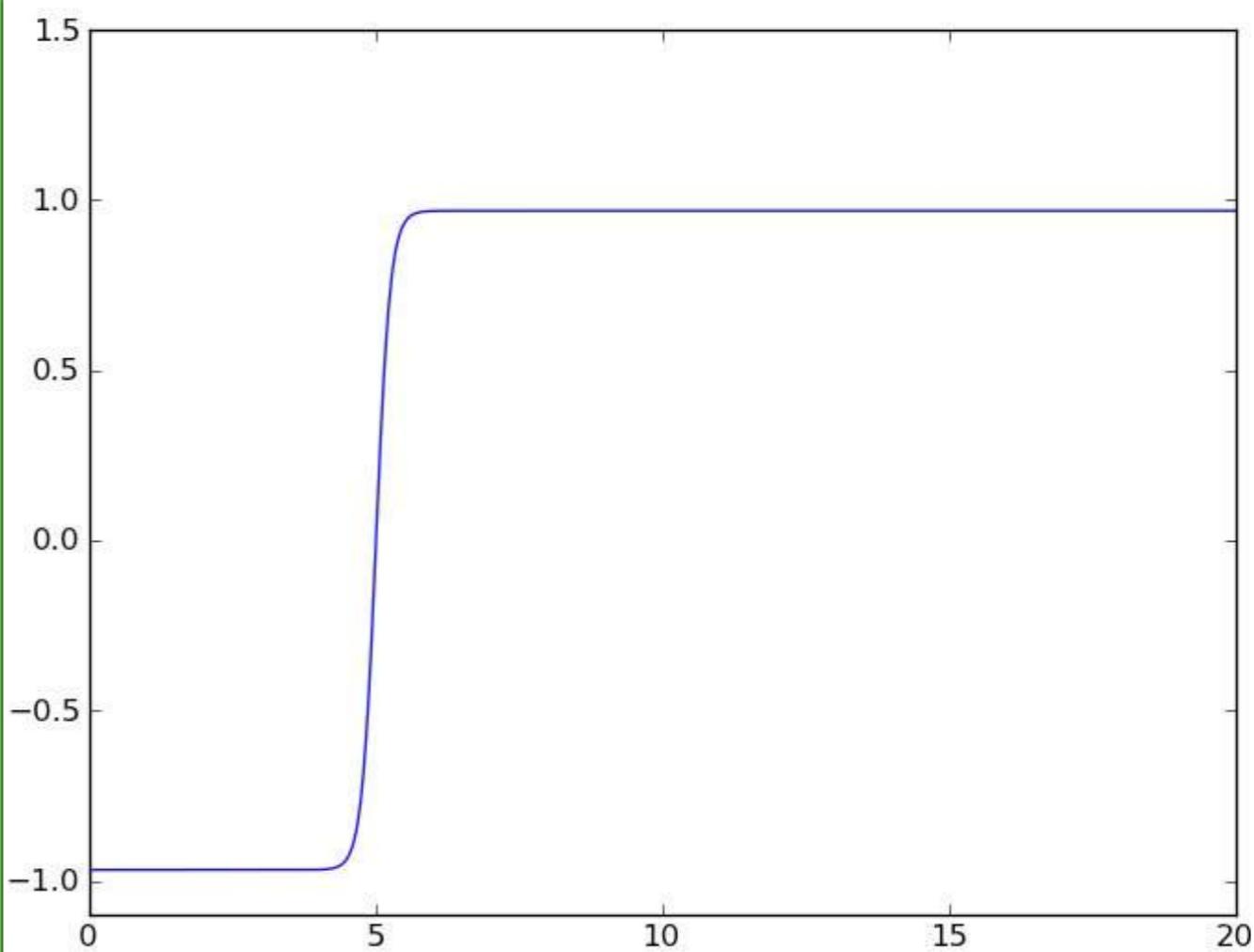
# Waves from Supernovae



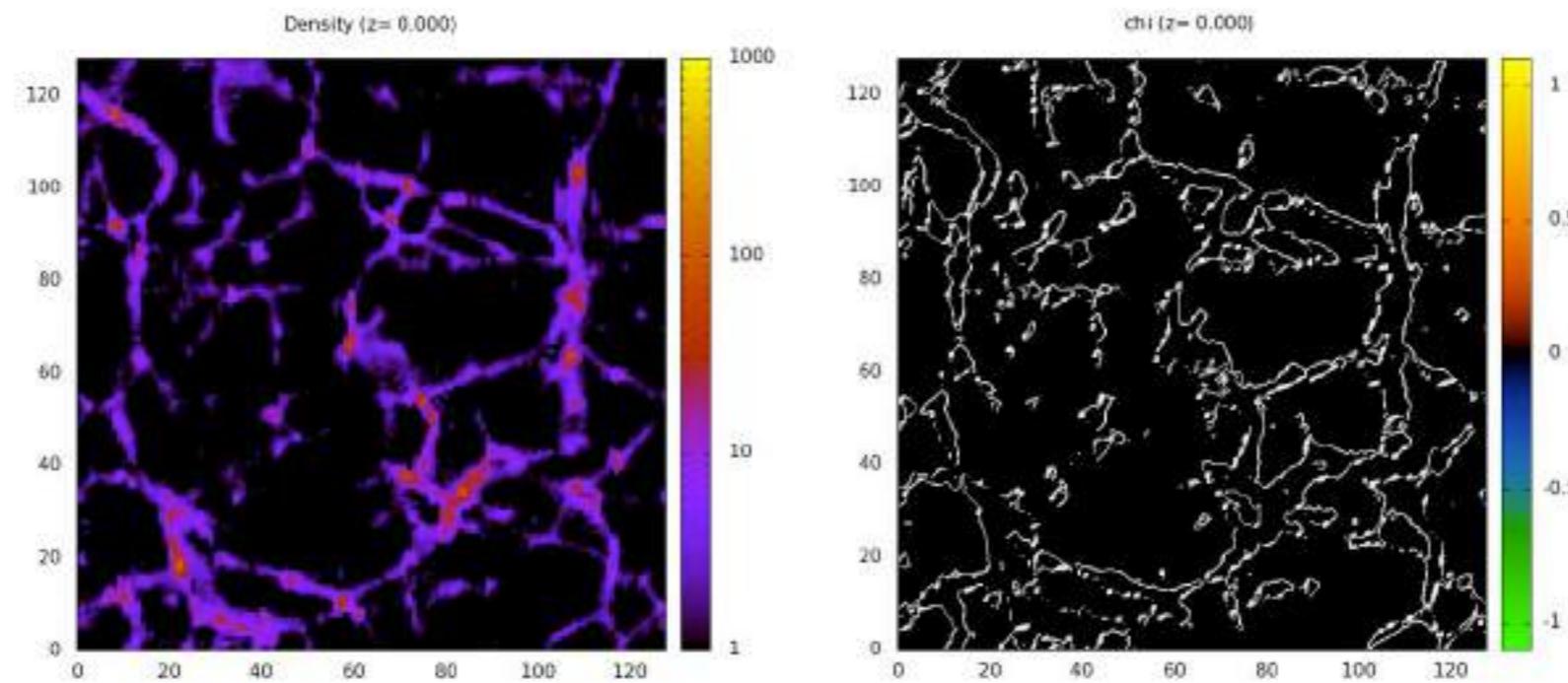
Chameleon Potential



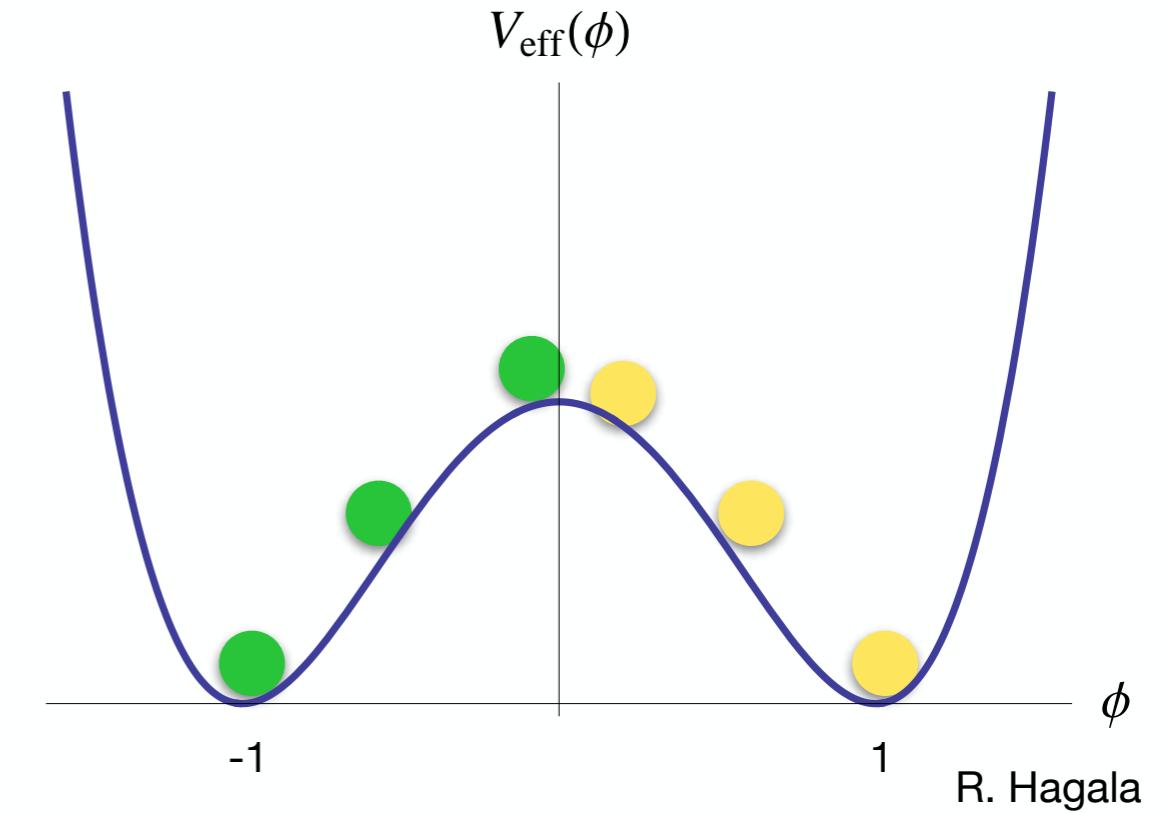
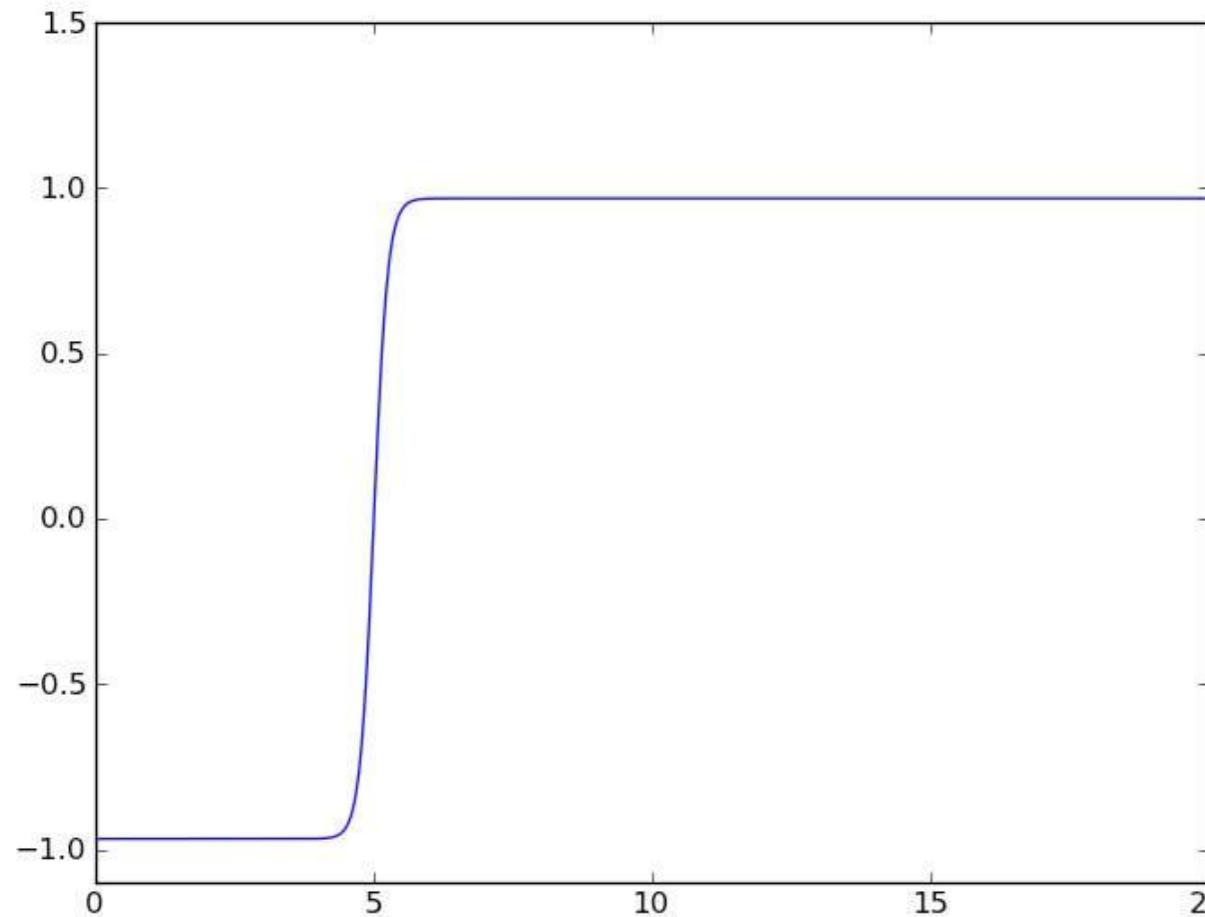
Profile of field



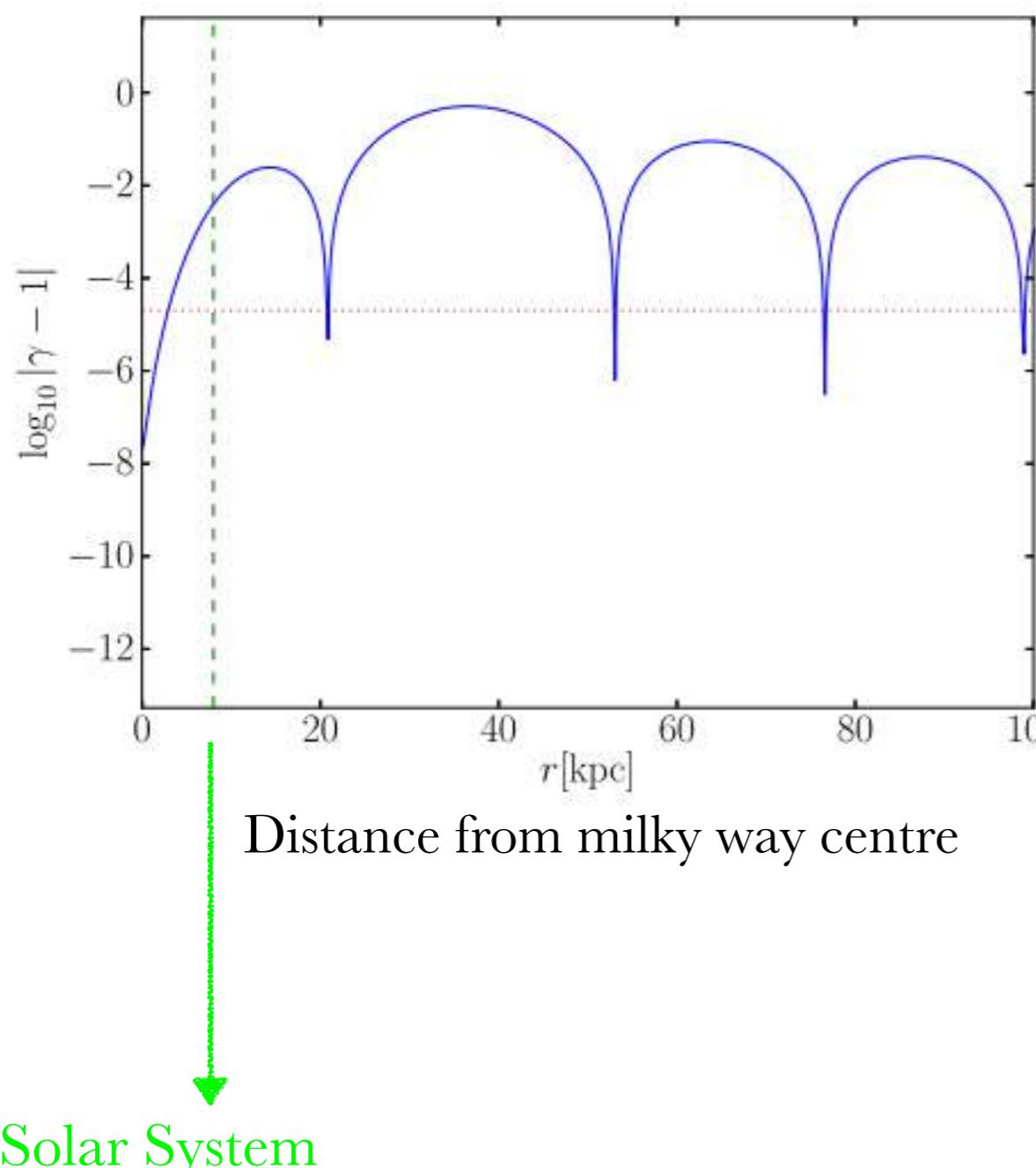
# Waves from collapse of domain walls



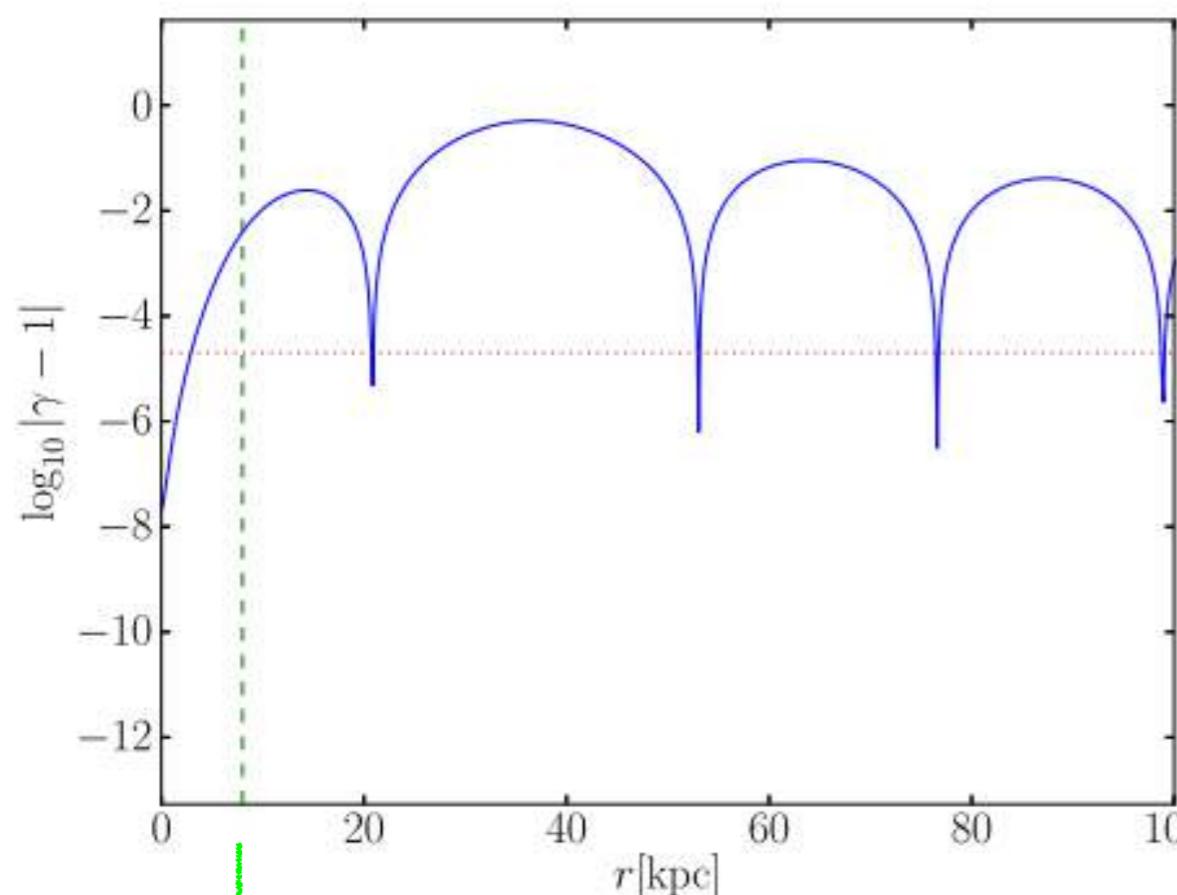
Profile of field



R. Hagala



$$\gamma - 1 = -\frac{\phi^2}{M^2} \frac{2}{\frac{\phi^2}{M^2} + 2\Psi(1 + \frac{\phi^2}{M^2})}$$



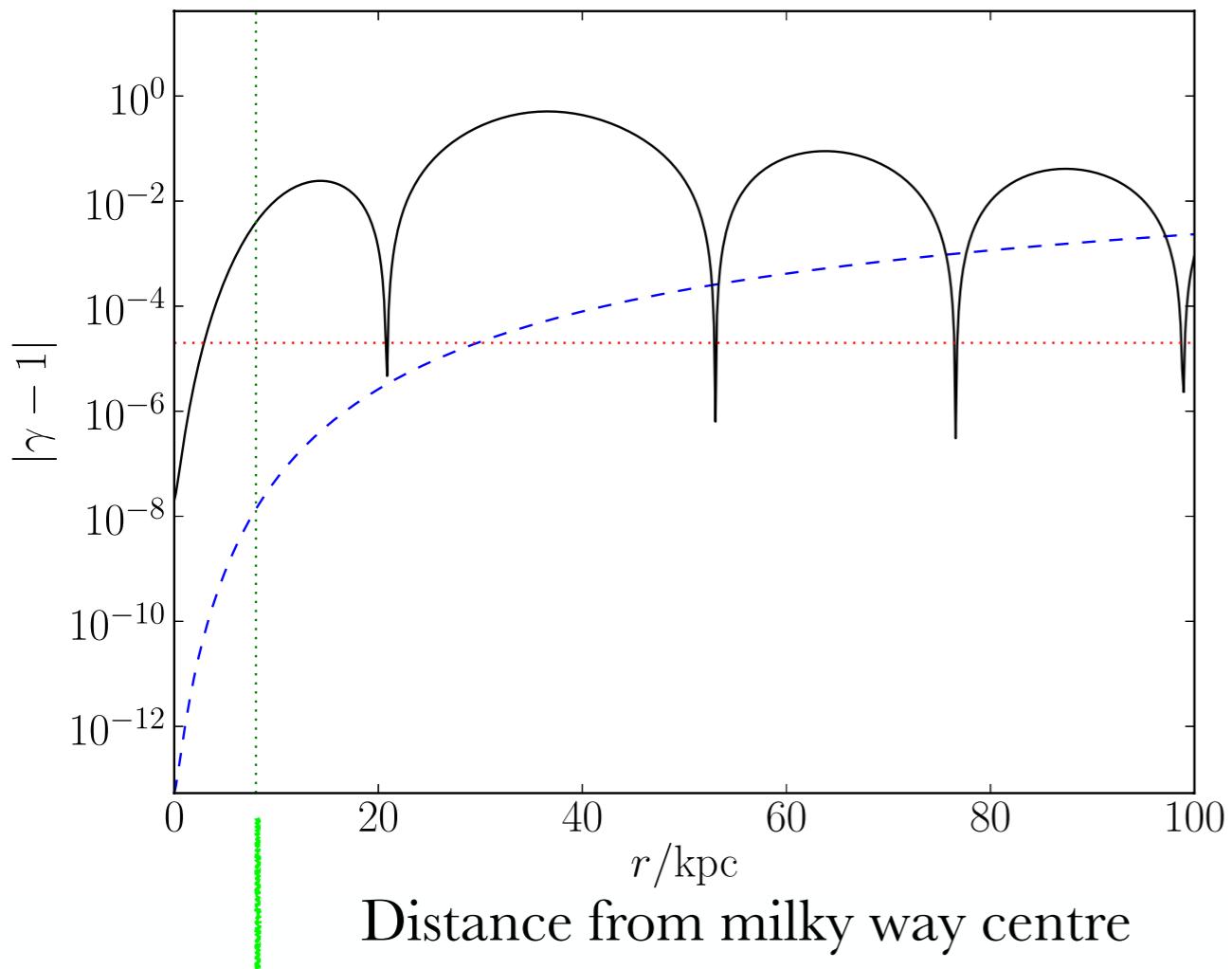
Distance from milky way centre

Solar System

Cassini bound



$$\gamma - 1 = -\frac{\phi^2}{M^2} \frac{2}{\frac{\phi^2}{M^2} + 2\Psi(1 + \frac{\phi^2}{M^2})}$$



Solar System

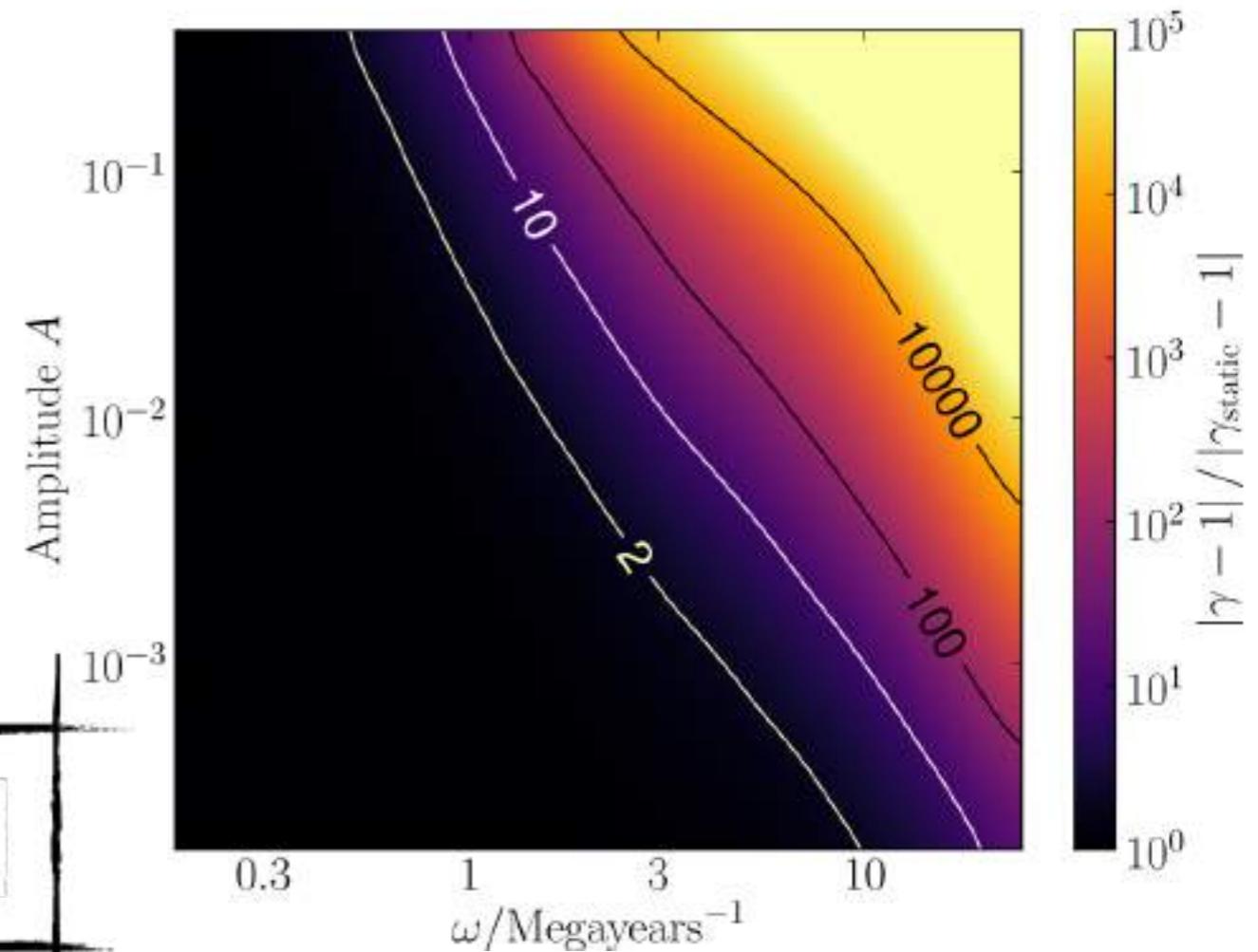
### Screening disrupted!

Scalar waves increase the PPN parameter and fifth force several orders in magnitude!

Oscillations are a smoking gun!

Cassini bound

$$\gamma - 1 = -\frac{\phi^2}{M^2} \frac{2}{\frac{\phi^2}{M^2} + 2\Psi(1 + \frac{\phi^2}{M^2})}$$



# Summary

- ▶ A light extra degree of freedom in the gravity sector is viable only if a screening mechanism is efficient to suppress it at local scales
- ▶ The viability and efficiency of screening mechanism generally relies on the quasi-static approximation
- ▶ Astrophysical events can create waves and the quasi-static approximation is no longer valid
- ▶ Waves diminish the screening mechanism efficiency in several orders of magnitude reducing the viability of many modified gravity theories